

Deconvoluting reversal modes in exchange-biased nanodots

Randy K. Dumas,¹ Chang-Peng Li,² Igor V. Roshchin,^{3,4} Ivan K. Schuller,² and Kai Liu^{1,*}

¹Physics Department, University of California, Davis, California 95616, USA

²Physics Department and Center for Advanced Nanoscience, University of California - San Diego, La Jolla, California 92093, USA

³Department of Physics and Astronomy, Texas A&M University, College Station, Texas, 77843, USA

⁴Materials Science and Engineering Program, Texas A&M University, College Station, Texas, 77843, USA

(Received 5 March 2012; revised manuscript received 18 June 2012; published 15 October 2012)

Ensemble-averaged exchange bias in arrays of Fe/FeF₂ nanodots has been deconvoluted into local, microscopic bias separately experienced by nanodots going through different reversal modes. The relative fraction of dots in each mode can be modified by exchange bias. Single-domain dots exhibit a simple loop shift, while vortex state dots have asymmetric shifts in the vortex nucleation and annihilation fields, manifesting local incomplete domain walls in these nanodots as magnetic vortices with tilted cores.

DOI: [10.1103/PhysRevB.86.144410](https://doi.org/10.1103/PhysRevB.86.144410)

PACS number(s): 75.30.Et, 75.60.-d, 75.70.Kw, 75.75.Fk

Ferromagnet/antiferromagnet (FM/AF) exchange-biased nanostructures¹ have critical applications in spin-valve type of spintronic devices.² They have also excited recent interests in antiferromagnetic spin-transfer-torque effect,^{3,4} low-power electrically controlled magnetic switching in multiferroics,⁵⁻⁷ and exchange-biased magnetic vortices,⁸⁻¹¹ which have novel potential applications in nanoelectronics and magnetic memory. Such nanostructures not only have structural heterogeneity across the FM/AF interface but also often have nanoscale confinements in the lateral dimension,¹²⁻¹⁵ leading to intriguing ground-state domain configurations and reversal mechanisms that are yet to be fully understood.

The exchange bias (EB) in these systems is conventionally obtained from the shift of the major hysteresis loop at the coercive fields, defined as H_E . However, some systems show very different unidirectional anisotropy energies when measured by reversible vs irreversible methods,^{16,17} and the anisotropy may be dependent upon sample *lateral* dimensions. For instance, in NiO/NiFe and FeMn/NiFe microstructures, the H_E is larger than in uniform films;¹⁸ a size dependence of EB is recently reported in Co/CoO nanostructures.¹⁹ These results suggest an H_E distribution in seemingly uniform films, raising the question whether the H_E obtained from the major loop (ML) is uniform across the entire sample or just a reflection of the weakest point. In prior studies of Fe/epitaxial-FeF₂ films, a distribution of H_E was shown to be centered on the ML value.²⁰ However, since the FM layer was continuous, variations of exchange coupling within the layer may mask the distribution of local EB. *A clear link is still missing between the ensemble-averaged H_E measured by magnetometry and the local, microscopic EB.* In addition, local EB in FM/AF hybrids can be tailored by nanopatterning.¹²⁻¹⁴ In particular, the local balance between magnetostatic and exchange energies in the FM can strongly modify the magnetization reversal mode, e.g., single-domain (SD) vs vortex state (VS) reversal. The interplay between EB and the VS reversal mode has led to a host of fascinating properties, including angular-dependent magnetization reversal,⁸ chirality control,¹⁰ and vortex imprinting into the AF.^{9,11}

Furthermore, depth-dependent magnetic configurations across the FM/AF interface are also important. Recent studies have highlighted the role of local incomplete domain walls

(LIDWs) in continuous films in generating the magnetization reversal asymmetry commonly observed.^{21,22} Such LIDWs form parallel to the FM/AF interface due to the competition between inhomogeneous interfacial exchange and the magnetic field, and their lateral extent can change within the film plane under field cycling. However, the existence and behavior of such LIDWs in EB nanostructures with confined lateral dimensions have remained unexplored.

In this paper we report the interplay between geometric confinement and EB in patterned arrays of Fe/FeF₂ nanodots of sizes close to the SD and VS reversal mode boundary for Fe. The ML EB has been deconvoluted into weighted averages of local EB separately experienced by SD and VS dots. The EB also modifies the relative fraction of dots that undergo VS and SD reversal modes. These findings establish the long missing link between macroscopic and local EB. Furthermore, we have observed unequal shifts in the bias manifested by the VS nucleation and annihilation fields, which provide first experimental demonstration of LIDWs in patterned structures as magnetic vortices with tilted cores.

Polycrystalline Fe (20 nm)/FeF₂ (50 nm)/Ag cap (5 nm) nanodots were grown on Si substrates using nanoporous alumina shadow masks in conjunction with electron beam evaporation and Ar ion etching.²³ The nanodot size is 67 ± 13 nm, and coexistence of SD and VS modes is expected from our prior studies on unbiased Fe dots.^{24,25} The dot center-to-center spacing is roughly twice the dot diameter, as illustrated in Fig. 1 inset, so that the dipolar interactions between the dots are negligible.²⁶

Magnetic properties were measured in a field applied in the plane of the nanodots using a vibrating sample magnetometer with a liquid helium flow cryostat. Major hysteresis loops were measured at room temperature (RT) and at 40 K, after two different cooling procedures. The field-cooling (FC) procedure involves cooling from RT through the FeF₂ Néel temperature ($T_N = 78$ K) in a 5 kOe applied field to 40 K. Subsequent measurements were then taken with the applied field parallel to the FC direction. The zero-field-cooling (ZFC) procedure involves ac demagnetizing the sample at RT and then cooling in zero applied field to 40 K. Note that no training effect was observed, and the field sweep rate (~ 12 Oe/s) was slow enough to avoid any relaxation processes.

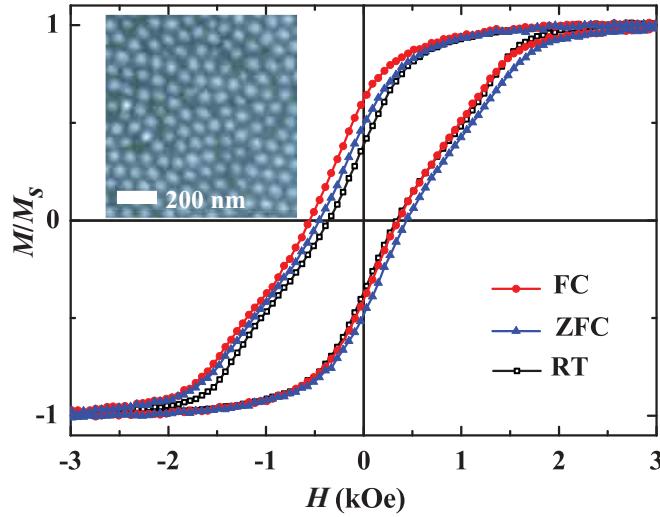


FIG. 1. (Color online) Major hysteresis loops measured at 40 K, after separate FC and ZFC procedures, in comparison with that measured at RT. Inset shows a scanning electron micrograph of 67-nm unbiased Fe nanodots.

The first-order reversal curve (FORC) technique^{24,25,27} was employed to investigate details of the magnetization reversal. The FORC distribution is defined as

$$\rho(H, H_R) \equiv -\frac{1}{2} \frac{\partial^2 M(H, H_R)/M_S}{\partial H \partial H_R}, \quad (1)$$

where H_R is the reversal field, H is the applied field, M is the magnetization, and M_S is the saturation magnetization. This distribution eliminates the purely reversible components of the magnetization. Thus, any nonzero ρ corresponds to irreversible switching processes. The FORC distribution can also be plotted in coordinates of (H_C, H_B) , where $H_C = (H - H_R)/2$ is the local coercive field and $H_B = (H + H_R)/2$ is the local interaction or bias field. The integration

$$I_{\text{irrev}} = \int \rho(H_C, H_B) dH_C dH_B \quad (2)$$

captures the total amount of irreversible magnetic switching.²⁰ A quantitative measure of the relative fraction of dots with a particular reversal mode is evaluated by selectively integrating the FORC distribution over the region of interest.

Magnetic hysteresis loops at 40 K are shown in Fig. 1. The FC sample displays a shifted loop and a coercivity enhancement from 348 Oe at RT to 456 Oe. Conventionally, H_E is determined from the ML shift according to

$$H_E^{\text{ML}} = (H_C^R + H_C^L)/2, \quad (3)$$

where H_C^R and H_C^L are the right and left coercive fields, respectively. Using this, an EB of $H_E^{\text{ML}} = -97$ Oe is found. Note that compared to Fe/epitaxial-FeF₂ films,²⁰ these nanodots exhibit an enhanced coercivity caused by the nanopatterning¹⁶ and a reduced EB due to the polycrystalline nature of the FeF₂.²⁸ In contrast, the ZFC sample shows no EB but an increased coercivity of 447 Oe at 40 K. The coercivity increase from RT to 40 K in the FC and ZFC samples is due to both the influence of the FeF₂ and the reduced temperature.²⁵

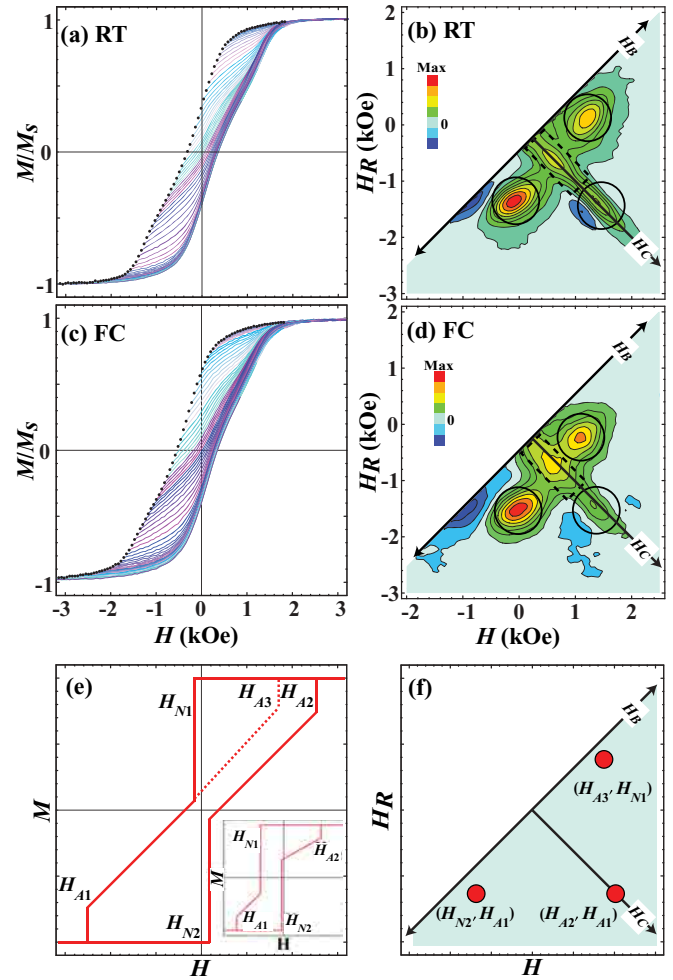


FIG. 2. (Color online) Families of FORCs, whose starting points are represented by black dots (a) at RT and (c) after FC to 40 K. The corresponding FORC distributions are shown as contour plots in (b) and (d), respectively. The circled regions in (b) and (d) indicate the highly irreversible processes associated with the VS nucleation and annihilation, while the dotted black box highlights the SD reversal mode. (e) Schematic hysteresis loop for unbiased vortices with nucleation field H_N and annihilation field H_A labeled. The inset shows an exchange-biased loop. (f) Schematic FORC distribution indicating the relationship between the three primary FORC peaks in the (H, H_R) coordinate system and the nucleation/annihilation fields shown in (e).

The RT FORCs and the corresponding FORC distribution in the (H, H_R) coordinate system are shown in Figs. 2(a) and 2(b), respectively. The FORC distribution is similar to that observed previously in comparably sized and unbiased Fe nanodot arrays, displaying a SD and VS reversal mode mixture.^{24,25,29} In Fig. 2(b) the FORC feature highlighted by the dashed box corresponds to the SD reversal mode; the other three peaks highlighted by circles correspond to the VS reversal mode. By projecting the former FORC distribution ρ along the bias field H_B axis, the EB of the SD dots can be determined. As shown in Fig. 3, for SD dots, the dM/dH_B peak at RT is centered on $H_B = 0$, consistent with an EB of $H_E^{\text{SD}} = 0$ Oe. In addition by selectively integrating the SD region of the FORC distribution and using Eq. (2), the fraction of dots reversing at RT by a SD reversal mode is $f_{\text{SD}} = 27 \pm 2\%$.³⁰

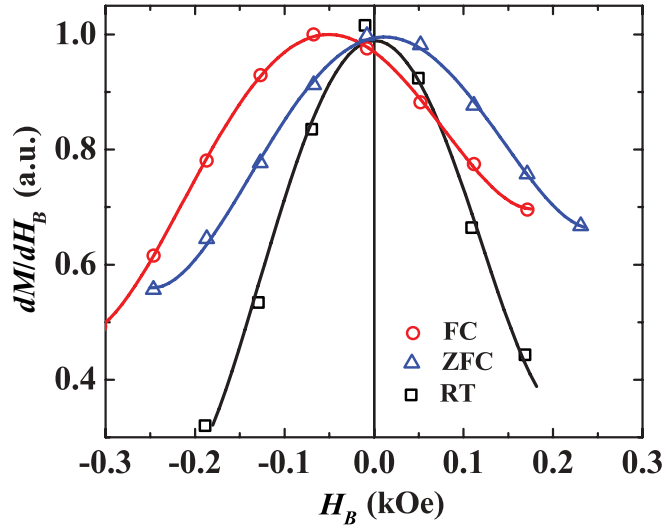


FIG. 3. (Color online) Projection of the FORC distribution along the H_B axis, dM/dH_B , for the SD reversal mode at RT, after FC, and after ZFC. The shift of the peak along the H_B axis indicates the exchange bias of the SD dots.

For the VS dots, FORC peaks and their widths are directly related to the VS nucleation/annihilation fields and their corresponding distributions. A schematic VS hysteresis loop is shown in Fig. 2(e). Along the descending (ascending) field branch, a vortex nucleates at H_{N1} (H_{N2}) and annihilates at H_{A1} (H_{A2}); as the field is reversed to positive saturation immediately after H_{N1} , the annihilation field is H_{A3} [dashed line in Fig. 2(e)]. Figure 2(f) is a schematic FORC distribution in the (H, H_R) coordinate system highlighting the relationship between the three primary FORC peaks and the nucleation/annihilation fields shown in Fig. 2(e). For the RT FORC distribution shown in Fig. 2(b), along the descending branch, the VS nucleation field ($H_{N1} = 140$ Oe) and annihilation field ($H_{A1} = -1340$ Oe) are symmetric to those along the ascending branch ($H_{N2} = -140$ Oe, $H_{A2} = 1340$ Oe).³¹ Thus, the dots that undergo VS reversal do not exhibit EB at RT either, as expected.

After FC the sample, the 40 K FORCs and the corresponding distribution are shown in Figs. 2(c) and 2(d), respectively. Although they appear qualitatively similar to those at RT, the peak positions and relative intensities in the FORC distribution are quite different. By integrating over the SD feature, we find the fraction of SD dots has increased to $f_{SD} = 32 \pm 3\%$. Additionally, as shown in Fig. 3, the SD peak is shifted toward negative-bias fields. The projection of $\rho(H_C, H_B)$ along the H_B axis, dM/dH_B , peaks at -50 Oe, indicating an EB of $H_E^{SD} = -50$ Oe. Interestingly, this value is nearly half the ML value, $H_E^{ML} = -97$ Oe, shown in Fig. 1. Note that H_E^{ML} effectively measures the shift in magnetization switching fields along the descending and ascending branches. Equivalently and analogous to Eq. (3), an EB can be calculated from the shift of the VS nucleation fields as follows:

$$H_E^{VS,N} = (H_{N1} + H_{N2})/2. \quad (4)$$

From the FORC distribution in Fig. 2(d), we find $H_{N1} = -265$ Oe and $H_{N2} = -13$ Oe; hence, $H_E^{VS,N} = -139$ Oe, significantly larger than the ML value. This key difference

illustrates that SD and VS dots are affected by EB differently for the same FC procedure. As the ML is an ensemble average of both SD and VS phases, we extract an “averaged” EB using the calculated EBs and phase fractions from the FORC analysis:

$$H_E^{FORC} = H_E^{SD} \times f_{SD} + H_E^{VS,N} \times (1 - f_{SD}). \quad (5)$$

This weight-averaged EB is -110 Oe, similar to the value found from the ML analysis. Thus, the conventionally determined H_E can be deconvoluted into distributions of local EB experienced by all the dots in the assembly.

Additionally, the FORC analysis of the VS reversal reveals a much more interesting and complex behavior than a simple shift of the ML would suggest. Analogous to Eqs. (3) and (4), an EB can be calculated from the shift of the VS annihilation fields:

$$H_E^{VS,A} = (H_{A1} + H_{A2})/2. \quad (6)$$

Using the VS peak positions in Fig. 2(d), we find $H_{A1} = -1523$ Oe and $H_{A2} = +1400$ Oe, thus $H_E^{VS,A} = -61$ Oe. This is significantly different from the ML bias, and much smaller than the shift of the nucleation fields $H_E^{VS,N}$ calculated using Eq. (4). If the VS loop was simply shifted along the applied field axis, the values calculated using Eqs. (4) and (6) should be identical, i.e., both H_N and H_A are displaced by the same amount. However, the FORC results indicate that the VS nucleation fields are experiencing a larger displacement than the annihilation fields. This results in an asymmetric VS ML, as schematically shown in the inset of Fig. 2(e), where the top portion of the loop is wider than the bottom.

Such an asymmetry is a signature of the LIDWs,^{21,22} which have spring-magnet-like³² varying magnetization along the depth of the FM, and are manifested in nanodots as magnetic vortices with tilted cores. The Monte Carlo simulations of Mejía-López *et al.*³³ model a nearly identical material system as the Fe/FeF₂ dots studied here. They find a pronounced asymmetric reversal of the VS ML due to a nonuniform magnetization through the thickness of the Fe layer. In fact, their simulated hysteresis loops are qualitatively similar to the Fig. 2(e) inset, which is representative of the H_N and H_A values found after FC the sample. The analytical calculations of Guslienko and Hoffmann³⁴ also confirm the type of asymmetry we observe. In their calculations the asymmetry is dictated by the quantity $\ln(R/R_C)/2$, where R is the radius of the dot and R_C the radius of the vortex core. Using values appropriate for this sample, $R = 33.5$ nm and $R_C = 8$ nm.³⁵

$$\ln(R/R_C)/2 < 1 \rightarrow |H_E^A| < |H_E^N|. \quad (7)$$

That is, the apparent shift of the nucleation fields is larger than that of the annihilation fields.³⁶ This happens because in relatively small dots, where the vortex core has an energy contribution comparable to that of the rest of the dot, the energy barrier for the vortex nucleation (which requires an area with out-of-plane magnetization) is much larger than the barrier for annihilation. Such a difference between the shifts of the nucleation and annihilation fields is indeed what we have observed experimentally. Just as in the Monte Carlo simulations, the observed asymmetry is a result of depth-dependent magnetization variations in the FM. In other

words, the degree of pinning induced by the AF is stronger near the FM/AF interface than at the free surface of the FM, leading to the formation of LIDWs and magnetization reversal asymmetry.²² The unidirectional magnetic anisotropy induced by the EB suppresses the VS. This suppression results in a delayed vortex nucleation and earlier vortex annihilation at the FM/AF interface during reversal, as compared to the free FM surface. As the lateral extent of such LIDWs is limited by the dot size and a vortex core exists in the dot, the noncollinear magnetic moments along the depth of the FM manifest as a tilted vortex core.

Finally, we examine the evolution of the SD phase fraction. As mentioned above, the SD fraction f_{SD} increases from $27 \pm 2\%$ at RT to $32 \pm 3\%$ in the FC sample at 40 K. As previously studied in unbiased Fe dots,²⁵ f_{SD} can simply increase at lower temperatures due to suppression of thermally assisted vortex nucleation/annihilation. In order to distinguish the influence of EB alone on the SD phase fraction, we have investigated nanodots at 40 K under ZFC. From the FORC analysis, we find that the projection of the SD reversal along the H_B axis (Fig. 3) and VS peaks show no EB. However, integration of the FORC diagram shows that after ZFC, f_{SD} is only $28 \pm 3\%$, which is smaller than that after FC and nearly the original RT value. After demagnetizing the sample at RT, most of the dots are in the VS. Upon ZFC to 40 K, while the lower temperature alone should increase f_{SD} , there is a competing mechanism that lowers f_{SD} . As shown by Sort *et al.*,⁹ the vortex structure can be imprinted into the AF upon ZFC, which then enhances and stabilizes the VS reversal. A lower f_{SD} is then expected

after ZFC relative to the FC procedure as we have observed. This relatively small change in f_{SD} indicates that it is only those SD dots nearest the SD-VS phase boundary that now reverse as VS entities after ZFC.

In summary, we have investigated the interplay between geometric confinement and EB in Fe/FeF₂ nanodots that contain a mixture of SD and VS reversal modes. The use of a FORC technique allows for the behavior of each reversal mode to be deconvoluted from the ML. After FC, the SD dots show a simple loop shift, whereas reversal of the VS dots is asymmetric, i.e., the nucleation fields exhibit a larger EB than the annihilation fields. These results agree with both Monte Carlo simulations and analytical calculations and confirm the existence of LIDWs along the depth of the FM and tilted vortex cores. The macroscopic EB extracted from conventional ML measurements is found to be a weighted average of the local EB experienced by the SD and VS nanodots. Moreover, the fraction of dots experiencing a given reversal mode depends on the cooling procedure, consistent with imprinting the vortex structure into the AF. These results unravel the crucial role played by local interfacial exchange on EB and demonstrate a technique to deconvolute the behavior of different reversal modes across a macroscopic array of nanomagnets.

This work has been supported by the NSF (Grant No. DMR-1008791 and No. ECCS-0925626), the AFOSR (Grant No. FA9550-10-1-0409), and the Texas A&M University-CONACyT collaborative research program.

*Corresponding author: kailiu@ucdavis.edu

- ¹J. Noguees, J. Sort, V. Langlais, V. Skumryev, S. Surinach, J. S. Munoz, and M. D. Baro, *Phys. Rep.* **422**, 65 (2005).
- ²B. Dieny, V. S. Speriosu, S. S. P. Parkin, B. A. Gurney, D. R. Wilhoit, and D. Mauri, *Phys. Rev. B* **43**, 1297 (1991).
- ³Z. Wei, A. Sharma, A. S. Nunez, P. M. Haney, R. A. Duine, J. Bass, A. H. MacDonald, and M. Tsoi, *Phys. Rev. Lett.* **98**, 116603 (2007).
- ⁴S. Urazhdin and N. Anthony, *Phys. Rev. Lett.* **99**, 046602 (2007).
- ⁵L. W. Martin, Y. H. Chu, M. B. Holcomb, M. Huijben, P. Yu, S. J. Han, D. Lee, S. X. Wang, and R. Ramesh, *Nano Lett.* **8**, 2050 (2008).
- ⁶J. T. Heron, M. Trassin, K. Ashraf, M. Gajek, Q. He, S. Y. Yang, D. E. Nikonov, Y. H. Chu, S. Salahuddin, and R. Ramesh, *Phys. Rev. Lett.* **107**, 217202 (2011).
- ⁷J. Allibe, S. Fusil, K. Bouzehouane, C. Daumont, D. Sando, E. Jacquet, C. Deranlot, M. Bibes, and A. Barthélémy, *Nano Lett.* **12**, 1141 (2012).
- ⁸J. Sort, A. Hoffmann, S. H. Chung, K. S. Buchanan, M. Grimsditch, M. D. Baro, B. Dieny, and J. Noguees, *Phys. Rev. Lett.* **95**, 067201 (2005).
- ⁹J. Sort, K. S. Buchanan, V. Novosad, A. Hoffmann, G. Salazar-Alvarez, A. Bollero, M. D. Baró, B. Dieny, and J. Nogués, *Phys. Rev. Lett.* **97**, 067201 (2006).
- ¹⁰W. Jung, F. J. Castaño, and C. A. Ross, *Phys. Rev. Lett.* **97**, 247209 (2006).
- ¹¹J. Wu, D. Carlton, J. S. Park, Y. Meng, E. Arenholz, A. Doran,

- A. T. Young, A. Scholl, C. Hwang, H. W. Zhao, J. Bokor, and Z. Q. Qiu, *Nat. Phys.* **7**, 303 (2011).
- ¹²K. Liu, S. M. Baker, M. Tuominen, T. P. Russell, and I. K. Schuller, *Phys. Rev. B* **63**, 060403 (2001).
- ¹³V. Baltz, J. Sort, S. Landis, B. Rodmacq, and B. Dieny, *Phys. Rev. Lett.* **94**, 117201 (2005).
- ¹⁴M. T. Rahman, N. N. Shams, D. S. Wang, and C. H. Lai, *Appl. Phys. Lett.* **94**, 082503 (2009).
- ¹⁵X. Sun, N. Frey Huls, A. Sigdel, and S. Sun, *Nano Lett.* **12**, 246 (2011).
- ¹⁶B. H. Miller and E. D. Dahlberg, *Appl. Phys. Lett.* **69**, 3932 (1996).
- ¹⁷V. Strom, B. Jonsson, K. Rao, and D. Dahlberg, *J. Appl. Phys.* **81**, 5003 (1997).
- ¹⁸J. G. Zhu, Y. F. Zheng, and X. D. Lin, *J. Appl. Phys.* **81**, 4336 (1997).
- ¹⁹S. Laureti, S. Y. Suck, H. Haas, E. Prestat, O. Bourgeois, and D. Givord, *Phys. Rev. Lett.* **108**, 077205 (2012).
- ²⁰J. Olamit, K. Liu, Z. P. Li, and I. K. Schuller, *Appl. Phys. Lett.* **90**, 032510 (2007).
- ²¹M. Kiwi, J. Mejia-Lopez, R. Portugal, and R. Ramirez, *Appl. Phys. Lett.* **75**, 3995 (1999).
- ²²Z. P. Li, O. Petravic, R. Morales, J. Olamit, X. Batlle, K. Liu, and I. K. Schuller, *Phys. Rev. Lett.* **96**, 217205 (2006).
- ²³C.-P. Li, I. V. Roshchin, X. Batlle, M. Viret, F. Ott, and I. K. Schuller, *J. Appl. Phys.* **100**, 074318 (2006).
- ²⁴R. K. Dumas, C.-P. Li, I. V. Roshchin, I. K. Schuller, and K. Liu, *Phys. Rev. B* **75**, 134405 (2007).

- ²⁵R. K. Dumas, K. Liu, C. P. Li, I. V. Roshchin, and I. K. Schuller, *Appl. Phys. Lett.* **91**, 202501 (2007).
- ²⁶J. Mejia-Lopez, D. Altbir, A. H. Romero, X. Battle, I. V. Roshchin, C. P. Li, and I. K. Schuller, *J. Appl. Phys.* **100**, 104319 (2006).
- ²⁷J. E. Davies, O. Hellwig, E. E. Fullerton, G. Denbeaux, J. B. Kortright, and K. Liu, *Phys. Rev. B* **70**, 224434 (2004).
- ²⁸J. Olamit, Z. P. Li, I. K. Schuller, and K. Liu, *Phys. Rev. B* **73**, 024413 (2006).
- ²⁹Note that the VS loops here have finite remanence, as discussed in Ref. 24, in contrast to completely pinched loops typically observed in larger dots.
- ³⁰The uncertainty is calculated by integrating over an area bound by ± 30 Oe, which accounts for the discrete nature of our experimental field spacing.
- ³¹We locally fit each peak with a fourth-order polynomial in order to accurately determine its location in field space. The uncertainty is approximately ± 5 Oe.
- ³²E. E. Fullerton, J. S. Jiang, M. Grimsditch, C. H. Sowers, and S. D. Bader, *Phys. Rev. B* **58**, 12193 (1998).
- ³³J. Mejia-Lopez, P. Soto, and D. Altbir, *Phys. Rev. B* **71**, 104422 (2005).
- ³⁴K. Y. Guslienko and A. Hoffmann, *Phys. Rev. Lett.* **97**, 107203 (2006).
- ³⁵I. V. Roshchin, C.-P. Li, H. Suhl, X. Battle, J. Mejia-Lopez, D. Altbir, A. H. Romero, S. Roy, S. K. Sinha, M. Fitzimmons, and I. K. Schuller, *Europhys. Lett.* **86**, 67008 (2009).
- ³⁶Using methods outlined in Ref. 34, the calculated $H_E^{V.S.N}$ and $H_E^{V.S.A}$ are, respectively, -330 Oe and -61 Oe, in qualitative agreement with experimental values.